

Hermite stuff

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1 Hermite stuff

We define the Hermite inner product by

$$\langle u, v \rangle_H = \int_{-\infty}^{\infty} u(x)v(x)e^{-x^2} dx.$$

The first four Hermite polynomials are

$$\begin{aligned}H_0(x) &= 1, \\H_1(x) &= 2x, \\H_2(x) &= 4x^2 - 2, \\H_3(x) &= 8x^3 - 12x.\end{aligned}$$

We can show that they are orthogonal to one another by computing the Hermite inner products.

First, note that by symmetry we have

$$\int_{-\infty}^{\infty} (\text{any odd function})e^{-x^2} dx = 0.$$

Thus automatically we find

$$\begin{aligned}\langle H_0, H_1 \rangle_H &= 0 & \langle H_0, H_3 \rangle_H &= 0 \\ \langle H_2, H_1 \rangle_H &= 0 & \langle H_2, H_3 \rangle_H &= 0\end{aligned}$$

For the rest of the integrals, we first note that

$$\int_{-\infty}^{\infty} e^{-x^2} dx = \sqrt{\pi}$$

In [28]: `integrate(e^(-x^2), x, -infinity, infinity)`

Out [28]: `sqrt(pi)`

Next we note that

$$\int_{-\infty}^{\infty} x^2 e^{-x^2} dx = \int_{-\infty}^{\infty} \underbrace{x}_f \underbrace{xe^{-x^2} dx}_{dg} = \underbrace{x}_f \underbrace{\left. \frac{-1}{2} e^{-x^2} dx \right|_{-\infty}^{\infty}}_g + \frac{1}{2} \int_{-\infty}^{\infty} e^{-x^2} dx = \frac{1}{2} \sqrt{\pi}$$

```
In [29]: integrate((x^2)*e^(-x^2),x,-infinity, infinity)
```

```
Out[29]: 1/2*sqrt(pi)
```

Similarly

$$\int_{-\infty}^{\infty} x^4 e^{-x^2} dx = \frac{3}{4}\sqrt{\pi}$$

```
In [30]: integrate((x^4)*e^(-x^2),x,-infinity, infinity)
```

```
Out[30]: 3/4*sqrt(pi)
```

Thus we compute

$$\begin{aligned}\langle H_0, H_2 \rangle &= \int_{-\infty}^{\infty} (1)(4x^2 - 2)dx = 4 \left(\frac{1}{2}\sqrt{\pi} \right) - 2(\sqrt{\pi}) = 0 \\ \langle H_1, H_3 \rangle &= \int_{-\infty}^{\infty} (2x)(8x^3 - 12x)dx = 16 \left(\frac{3}{4}\sqrt{\pi} \right) - 24 \left(\frac{1}{2}\sqrt{\pi} \right) = 0\end{aligned}$$

Of course we could have had the computer do this for us:

```
In [31]: # define Hermite polynomials and the Hermite weight function
```

```
var('x')
H0(x)=1
H1(x) = 2*x
H2(x) = 4*x^2 -2
H3(x) = 8*x^3 - 12*x
w(x) = e^(-x^2)
```

```
In [32]: # compute inner product
```

```
integrate( H1(x)*H3(x)*w(x) ,x,-infinity, infinity)
```

```
Out[32]: 0
```

We can also compute the norms of the Hermite polynomials. We do this with the identity $\|u\|_H^2 = \langle u, u \rangle_H$.

$$\begin{aligned}\|H_0\|_H^2 &= \sqrt{\pi} \\ \|H_1\|_H^2 &= 2\sqrt{\pi} \\ \|H_2\|_H^2 &= 8\sqrt{\pi} \\ \|H_3\|_H^2 &= 48\sqrt{\pi}\end{aligned}$$

```
In [33]: latex(integrate( H3(x)*H3(x)*w(x) ,x,-infinity, infinity))
```

```
Out[33]: 48 \, \sqrt{\pi}
```

And yes, I did indeed have Sage give me the LaTeX code directly. Cause that's a thing.

2 Approximating $f(x) = e^x$

Let's consider the function $f(x) = e^x$. First we show that $\|f\|_H$ is finite. To do this by hand involves completing the square:

$$\langle f, f \rangle_H = \int_{-\infty}^{\infty} e^{2x-x^2} dx = \int_{-\infty}^{\infty} e^{-(x-1)^2+1} dx = e \int_{-\infty}^{\infty} e^{-(x-1)^2} dx = e\sqrt{\pi}$$

Let's check this with the machine:

```
In [34]: f(x) = e^x
         latex(integrate( f(x)*f(x)*w(x) ,x,-infinity, infinity))
```

```
Out[34]: \sqrt{\pi} e
```

The machine agrees, although I take exception to the decision to put $\sqrt{\pi}$ before e . Now we compute the coefficients

$$\alpha_k = \frac{\langle H_k, f \rangle_H}{\|H_k\|_H^2}$$

Using the machine, I find:

$$\begin{aligned}\alpha_0 &= e^{\frac{1}{4}} \\ \alpha_1 &= \frac{1}{2} e^{\frac{1}{4}} \\ \alpha_2 &= \frac{1}{8} e^{\frac{1}{4}} \\ \alpha_3 &= \frac{1}{48} e^{\frac{1}{4}}\end{aligned}$$

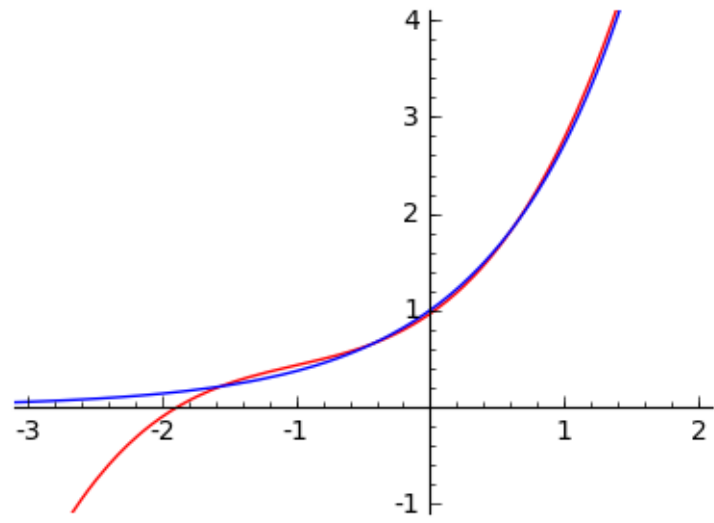
```
In [38]: H(x) = H3(x)
         latex(integrate( H(x)*f(x)*w(x) ,x,-infinity, infinity)/integrate( H(x)*H(x)*w(x) ,x,-i
```

```
Out[38]: \frac{1}{48} \, e^{\frac{1}{4}}
```

Here is a plot of the function f in blue and the approximation in red. Not bad!

```
In [53]: fH(x) = (e^(1/4))*H0(x) + (1/2*e^(1/4))*H1(x) + (1/8*e^(1/4))*H2(x) + (1/48*e^(1/4))*H3(x)
```

```
fHplot = plot(fH(x),(x,-10,10), color = 'red')
fplot = plot(f(x), (x,-10, 10))
mainplot = fHplot + fplot
mainplot.show(xmin=-3, xmax = 2, ymin = -1,ymax = 4,figsize = [4,3])
```



In []: